

3.35 A spin $\frac{1}{2}$ particle is in an orbital angular momentum $l = 1$ state.

- Starting with the total angular momentum state $|j, m\rangle = |j, j\rangle$ for $j = 3/2$, use the operator $\mathbf{J} = \mathbf{L} + \mathbf{S}$ to construct all the states $|j, m_j\rangle$ in terms of the eigenstates $|l, m_l\rangle$ for $\mathbf{L}^2$ and L_z , and the eigenstates $|1/2, \pm 1/2\rangle$ for $\mathbf{S}^2$ and S_z . Check your answers against any available table of Clebsch-Gordan coefficients; see Appendix E.
- If the particle is in a total angular-momentum eigenstate with z -component $+\hbar/2$, calculate the probability of finding the z -component of the spin of the particle to have the value $m_s = +1/2$.

Appendix

44. Clebsch-Gordan Coefficients, Spherical Harmonics, and d Functions

Note: A square-root sign is to be understood over *every* coefficient, e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation:

J	J	...
M	M	...
m_1	m_2	
m_1	m_2	Coefficients
$\vdots$	$\vdots$	
$\vdots$	$\vdots$	

$$1/2 \times 1/2$$

	1		
	+1	1	0
+1/2 +1/2	1	0	0
+1/2 -1/2	1/2	1/2	1
-1/2 +1/2	1/2	-1/2	-1
-1/2 -1/2	1		

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$$

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$$

$$Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$$

$$Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$$

$$Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$$

$$3/2 \times 1$$

	5/2	
	+5/2	5/2 3/2

$$1 \times 1/2$$

	3/2		
	+3/2	3/2 1/2	
+1 +1/2	1	+1/2 +1/2	
+1 -1/2	1/3	2/3	3/2 1/2
0 +1/2	2/3	-1/3	-1/2 -1/2
0 -1/2	2/3	1/3	3/2
-1 -1/2	1/3	-2/3	-3/2
2 \times 1	3		
	+3	3 2	
-1 -1/2	1		

$$3/2 \times 1/2$$

	2		
	+2	2 1	
+3/2 +1/2	1	+1 +1	
+3/2 -1/2	1/4	3/4	2 1
+1/2 +1/2	3/4	-1/4	0 0
+1/2 -1/2	1/2	1/2	2 1

a. 풀이

L_z 에 대한 eigen value 가 $\hbar/2$ ($m = \frac{3}{2}$) 인 상태는

$\{J^2, J_z\}$ basis 와 $\{L^2, L_z, S^2, S_z\}$ basis 모두에서 eigen ket 이다. 따라서,

$$|j = \frac{3}{2}, m = \frac{3}{2}\rangle = |m_l = 1, m_s = \frac{1}{2}\rangle$$

이 상태에서 lowering operator $J_- = L_- + S_-$ 를 적용.

$$J_- |j = \frac{3}{2}, m = \frac{3}{2}\rangle = \sqrt{(\frac{3}{2} + \frac{3}{2})(+1)} |j = \frac{3}{2}, m = \frac{1}{2}\rangle = \sqrt{3} |j = \frac{3}{2}, m = \frac{1}{2}\rangle$$

$$(L_- + S_-) |m_l = 1, m_s = \frac{1}{2}\rangle = \sqrt{2} |m_l = 0, m_s = \frac{1}{2}\rangle + |m_l = 1, m_s = -\frac{1}{2}\rangle$$

$$|j=\frac{3}{2}, m=\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |m_l=0, m_s=\frac{1}{2}\rangle + \frac{1}{\sqrt{3}} |m_l=1, m_s=-\frac{1}{2}\rangle$$

한 번 더 양변에 $J_- = L_- + S_-$ 를 적용.

$$\begin{aligned} 2 |j=\frac{3}{2}, m=-\frac{1}{2}\rangle &= \sqrt{\frac{2}{3}} \left\{ \sqrt{2} |m_l=-1, m_s=\frac{1}{2}\rangle + |m_l=0, m_s=-\frac{1}{2}\rangle \right. \\ &\quad \left. + \frac{1}{\sqrt{3}} \cdot \sqrt{2} |m_l=0, m_s=-\frac{1}{2}\rangle \right\} \\ &= \frac{2}{\sqrt{3}} |m_l=-1, m_s=\frac{1}{2}\rangle + 2 \cdot \sqrt{\frac{2}{3}} |m_l=0, m_s=-\frac{1}{2}\rangle \end{aligned}$$

$$|j=\frac{3}{2}, m=-\frac{1}{2}\rangle = \frac{1}{\sqrt{3}} |m_l=-1, m_s=\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |m_l=0, m_s=-\frac{1}{2}\rangle$$

다시 적용. ... 하기엔, 마지막 결과는 변하다.

$$|j=\frac{3}{2}, m=-\frac{3}{2}\rangle = |m_l=-1, m_s=-\frac{1}{2}\rangle$$

우리가 아직 모르는 것은,

$$\begin{cases} |j=\frac{1}{2}, m=\frac{1}{2}\rangle = a |m_l=1, m_s=-\frac{1}{2}\rangle + b |m_l=0, m_s=\frac{1}{2}\rangle \\ |j=\frac{1}{2}, m=-\frac{1}{2}\rangle = c |m_l=0, m_s=-\frac{1}{2}\rangle + d |m_l=-1, m_s=\frac{1}{2}\rangle \end{cases}$$

이 관계에서 계수 a, b, c, d 를 모른다. 이걸 그냥 orthogonality 이용해서 구할 거다.

$$\begin{cases} |j=\frac{3}{2}, m=\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |m_l=0, m_s=\frac{1}{2}\rangle + \frac{1}{\sqrt{3}} |m_l=1, m_s=-\frac{1}{2}\rangle \\ |j=\frac{1}{2}, m=\frac{1}{2}\rangle = a |m_l=1, m_s=-\frac{1}{2}\rangle + b |m_l=0, m_s=\frac{1}{2}\rangle \end{cases}$$

← 이미 구한 거

$$\langle j=\frac{1}{2}, m=\frac{1}{2} | j=\frac{3}{2}, m=\frac{1}{2} \rangle = \frac{a}{\sqrt{3}} + b \sqrt{\frac{2}{3}} = 0$$

$$a = -\sqrt{2} b.$$

이 정보만으로는 a 와 b 의 부호를 결정할 수 없다.

convention에 따라, m_l 이 큰 $|m_l=1, m_s=-\frac{1}{2}\rangle$ 의 coefficient를 양수로 잡는다.

$|a|^2 + |b|^2 = 1$ normalization 을 진행하면,

$$a = \sqrt{\frac{2}{3}} \quad b = -\frac{1}{\sqrt{3}}$$

$$\begin{cases} |j=\frac{3}{2}, m=-\frac{1}{2}\rangle = \frac{1}{\sqrt{3}} |m_l=-1, m_s=\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |m_l=0, m_s=-\frac{1}{2}\rangle \\ |j=\frac{1}{2}, m=-\frac{1}{2}\rangle = c |m_l=0, m_s=-\frac{1}{2}\rangle + d |m_l=-1, m_s=\frac{1}{2}\rangle \end{cases}$$

이미 구한 거.

$$\langle j=\frac{1}{2}, m=-\frac{1}{2} | j=\frac{3}{2}, m=-\frac{1}{2} \rangle = c \sqrt{\frac{2}{3}} + d \frac{1}{\sqrt{3}} = 0$$

$d = -\sqrt{2}c$, 역시 convention에 따라 $c > 0$ 으로 잡고, normalization,

$$c = \frac{1}{\sqrt{3}}, \quad d = -\sqrt{\frac{2}{3}}$$

결과를 모두 모아 보면,

1 × 1/2		3/2			
		+3/2	3/2	1/2	
+1	+1/2	1	+1/2	+1/2	
		+1 -1/2	1/3	2/3	3/2
		0 +1/2	2/3	-1/3	1/2
			0 -1/2	2/3	1/3
			-1 +1/2	1/3	-2/3
				-1 -1/2	3/2
2 × 1	3				
	+3	3	2		
					1

$$\begin{cases} |j=\frac{3}{2}, m=\frac{3}{2}\rangle = |m_l=1, m_s=\frac{1}{2}\rangle \\ |j=\frac{3}{2}, m=\frac{1}{2}\rangle = \frac{1}{\sqrt{3}} |m_l=1, m_s=-\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |m_l=0, m_s=\frac{1}{2}\rangle \\ |j=\frac{1}{2}, m=\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |m_l=1, m_s=-\frac{1}{2}\rangle - \frac{1}{\sqrt{3}} |m_l=0, m_s=\frac{1}{2}\rangle \\ |j=\frac{3}{2}, m=-\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |m_l=0, m_s=-\frac{1}{2}\rangle + \frac{1}{\sqrt{3}} |m_l=-1, m_s=\frac{1}{2}\rangle \\ |j=\frac{1}{2}, m=-\frac{1}{2}\rangle = \frac{1}{\sqrt{3}} |m_l=0, m_s=-\frac{1}{2}\rangle - \sqrt{\frac{2}{3}} |m_l=-1, m_s=\frac{1}{2}\rangle \\ |j=\frac{3}{2}, m=-\frac{3}{2}\rangle = |m_l=-1, m_s=-\frac{1}{2}\rangle \end{cases}$$

결과가 기존에 알려진 CG와 일치한다!

b. 풀이 상태의 j 값이 명확하지 않으니 두 경우로 나누어 고자.

$$|j=\frac{3}{2}, m=\frac{1}{2}\rangle \text{인 경우, } m_s=\frac{1}{2} \text{ 일 확률은 } |\langle m_l=0, m_s=\frac{1}{2} | j=\frac{3}{2}, m=\frac{1}{2} \rangle|^2 = \frac{2}{3}$$

$$|j=\frac{1}{2}, m=\frac{1}{2}\rangle \text{인 경우, } m_s=\frac{1}{2} \text{ 일 확률은 } |\langle m_l=0, m_s=\frac{1}{2} | j=\frac{1}{2}, m=\frac{1}{2} \rangle|^2 = \frac{1}{3}$$

3.42 Consider a spherical tensor of rank 1 (that is, a vector)

$$V_{\pm 1}^{(1)} = \mp \frac{V_x \pm iV_y}{\sqrt{2}}, \quad V_0^{(1)} = V_z.$$

Using the expression for $d^{(j=1)}$ given in Problem 3.38, evaluate

$$\sum_{q'} d_{qq'}^{(1)}(\beta) V_{q'}^{(1)}$$

and show that your results are just what you expect from the transformation properties of $V_{x,y,z}$ under rotations about the y -axis.

problem 3.38

$$d^{(j=1)}(\beta) = \begin{pmatrix} \frac{1}{2}(1+\cos\beta) & (-\frac{1}{\sqrt{2}})\sin\beta & \frac{1}{2}(1-\cos\beta) \\ \frac{1}{\sqrt{2}}\sin\beta & \cos\beta & -\frac{1}{\sqrt{2}}\sin\beta \\ \frac{1}{2}(1-\cos\beta) & \frac{1}{\sqrt{2}}\sin\beta & \frac{1}{2}(1+\cos\beta) \end{pmatrix}$$

$$\sum_{q'} d_{qq'}^{(1)} V_{q'}^{(1)} = \begin{pmatrix} \frac{1}{2}(1+\cos\beta) & (-\frac{1}{\sqrt{2}})\sin\beta & \frac{1}{2}(1-\cos\beta) \\ \frac{1}{\sqrt{2}}\sin\beta & \cos\beta & -\frac{1}{\sqrt{2}}\sin\beta \\ \frac{1}{2}(1-\cos\beta) & \frac{1}{\sqrt{2}}\sin\beta & \frac{1}{2}(1+\cos\beta) \end{pmatrix} \begin{pmatrix} \frac{-V_x - iV_y}{\sqrt{2}} \\ V_z \\ \frac{V_x - iV_y}{\sqrt{2}} \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1}{\sqrt{2}}(-\cos\beta V_x - iV_y - \sin\beta V_z) \\ -\sin\beta V_x + \cos\beta V_z \\ \frac{1}{\sqrt{2}}(\cos\beta V_x - iV_y + \sin\beta V_z) \end{pmatrix} = \begin{pmatrix} \frac{-V'_x - iV'_y}{\sqrt{2}} \\ V'_z \\ \frac{V'_x - iV'_y}{\sqrt{2}} \end{pmatrix}$$

$$R_y(\beta) = \begin{pmatrix} \cos\beta & 0 & \sin\beta \\ 0 & 1 & 0 \\ -\sin\beta & 0 & \cos\beta \end{pmatrix}$$

$$\begin{pmatrix} V'_x \\ V'_y \\ V'_z \end{pmatrix} = \begin{pmatrix} \cos\beta & 0 & \sin\beta \\ 0 & 1 & 0 \\ -\sin\beta & 0 & \cos\beta \end{pmatrix} \begin{pmatrix} V_x \\ V_y \\ V_z \end{pmatrix} = \begin{pmatrix} \cos\beta V_x + \sin\beta V_z \\ V_y \\ -\sin\beta V_x + \cos\beta V_z \end{pmatrix}$$

$$\begin{cases} \underline{V'_x} + i V'_y = \underline{\cos\beta V_x + \sin\beta V_z} + i V_y \\ V'_z = -\sin\beta V_x + \cos\beta V_z \\ \underline{V'_x} - i V'_y = \underline{\cos\beta V_x + \sin\beta V_z} - i V_y \end{cases}$$

$V_{x,y,z}$ 가 다음 변환식을 따른다는 걸 알수있다.

$$\begin{cases} V'_x = \cos\beta V_x + \sin\beta V_z \\ V'_y = V_y \\ V'_z = -\sin\beta V_x + \cos\beta V_z \end{cases}$$